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Simulations of transport in porous media



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Simulations of transport in porous media

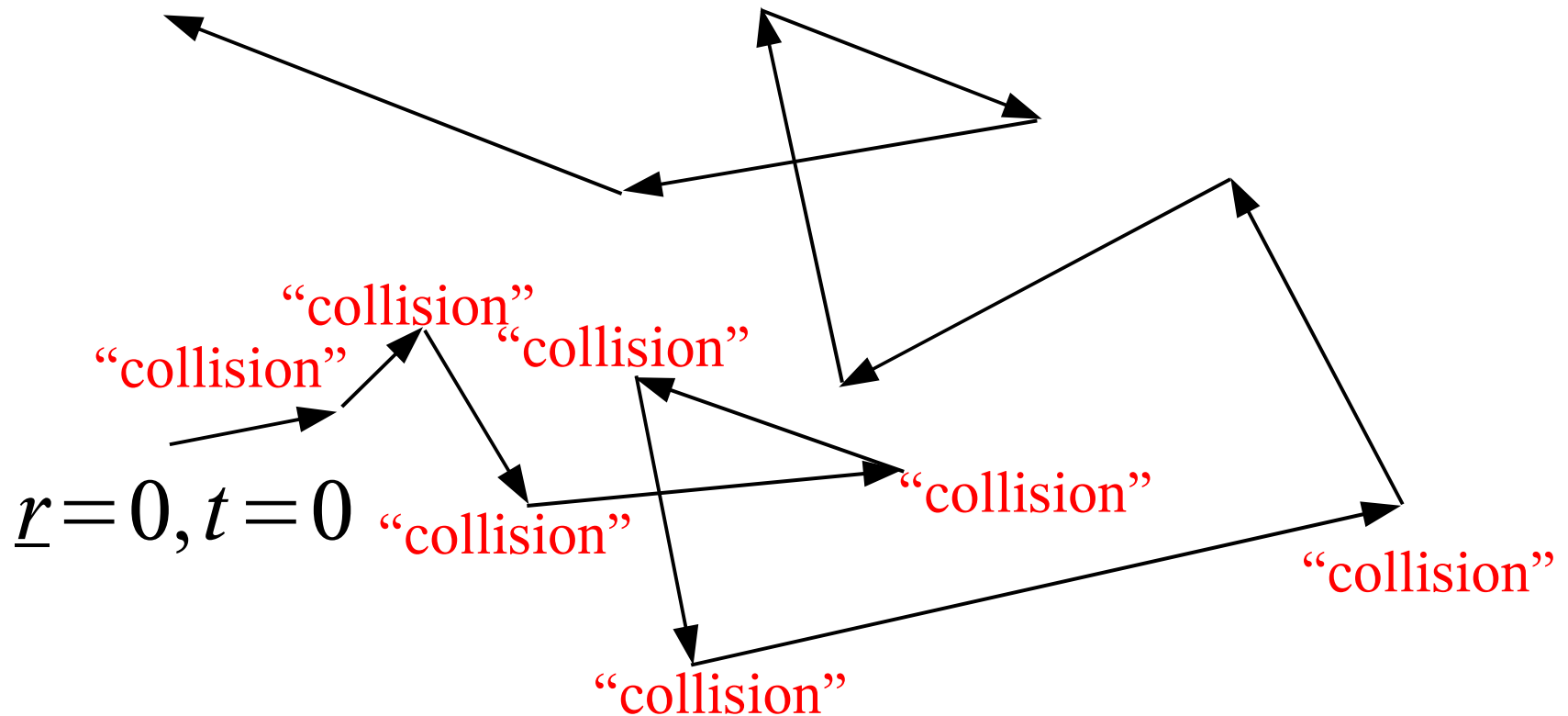
1. Diffusion as continuous time random walk
2. Types of diffusion
3. Creation of the model porous structures
4. Diffusion in model porous structures



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Diffusion as continuous time random walk

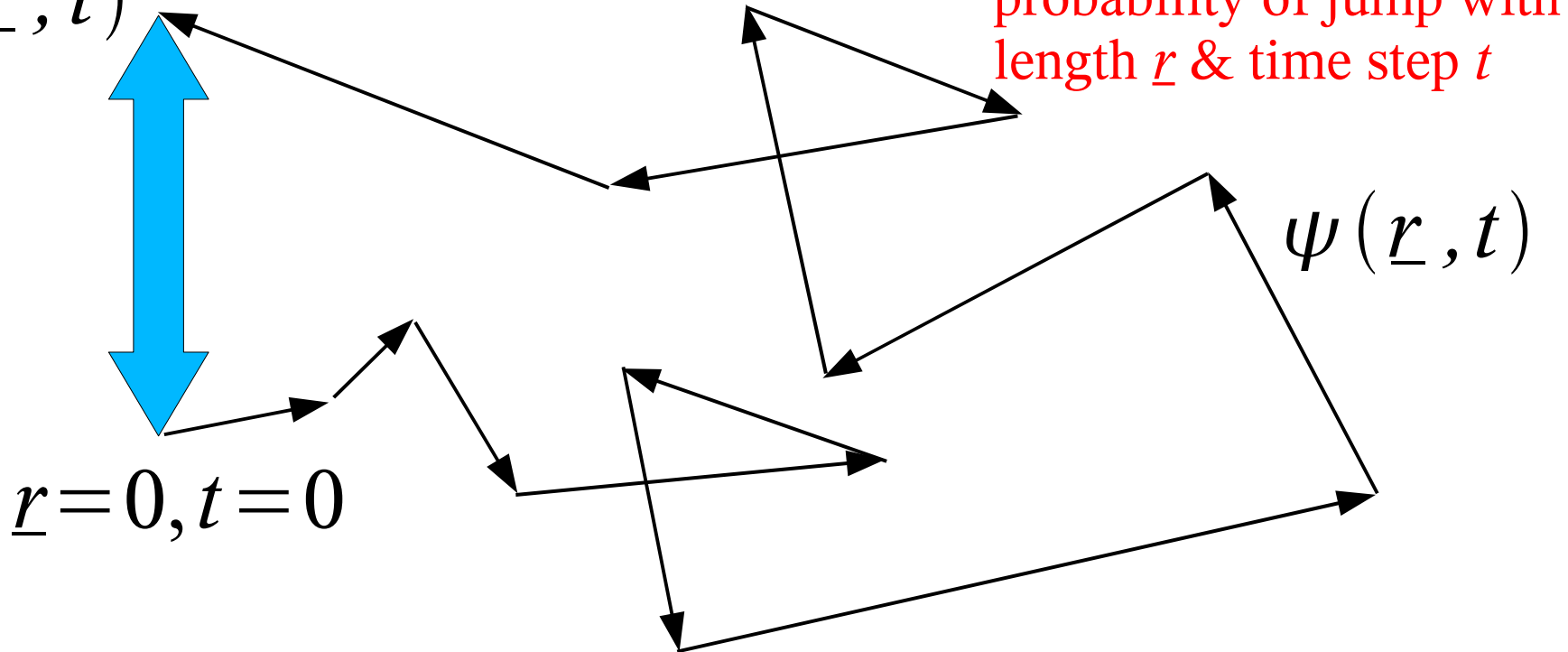


with: (1) another particle, **(2) porous structure walls,**
(3) crystal surface, (4) bubble surface, etc.

Diffusion as continuous time random walk

probability that particle is found at position \underline{r} at time t

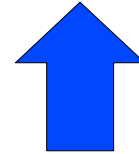
$p(\underline{r}, t)$



$$\langle r^2(t) \rangle = \int_{R_D} d\underline{r} r^2 p(\underline{r}, t) \quad \text{mean square displacement}$$

Diffusion as continuous time random walk

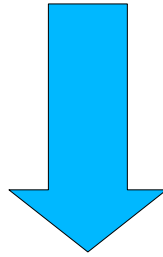
$$\frac{\partial}{\partial t} p(\underline{r}; t) = \int_{R_D} d\underline{r}' \int_0^{\infty} dt' \underbrace{K(\underline{r} - \underline{r}', t - t')}_{\text{kernel}} \underbrace{p(\underline{r}', t')}_{\text{probability of previous position}}$$



What is form of kernel?

Diffusion as continuous time random walk

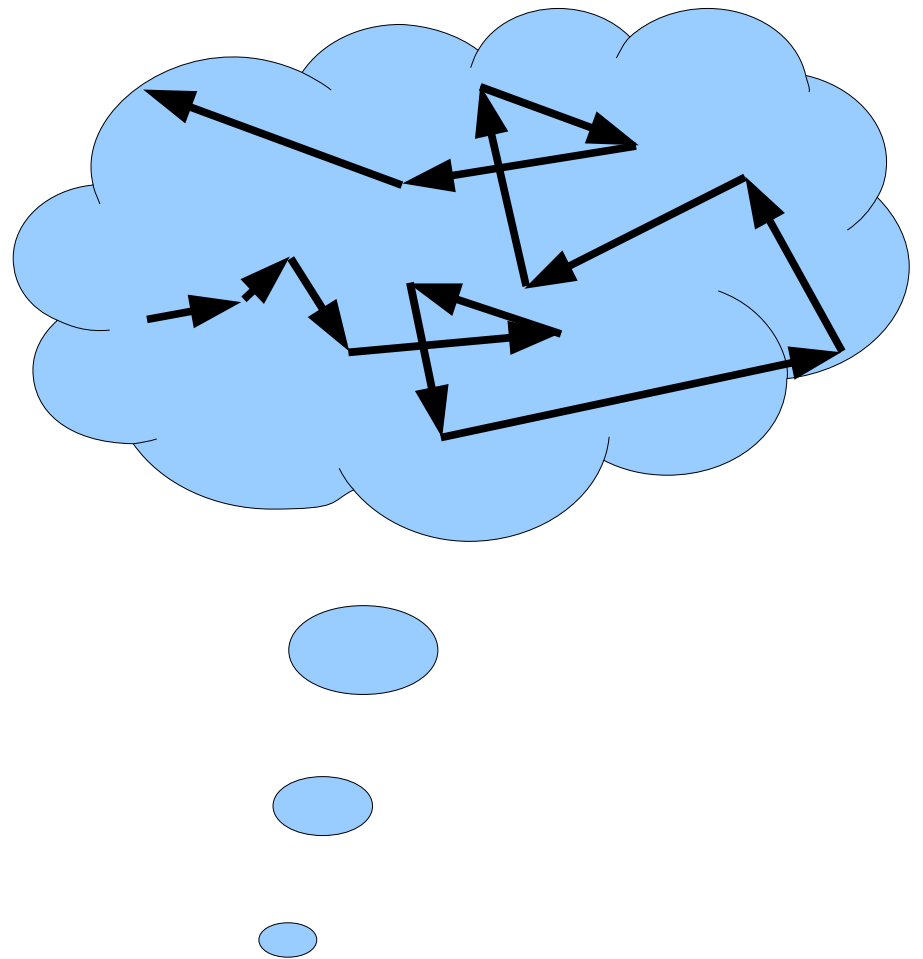
$$\frac{\partial}{\partial t} p(\underline{r}; t) = \int_{R_D} d\underline{r}' \int_0^{\infty} dt' K(\underline{r} - \underline{r}', t - t') p(\underline{r}', t')$$



$$ip(\underline{k}; u) - 1 = K(\underline{k}; u) p(\underline{k}; u)$$

Diffusion as continuous time random walk

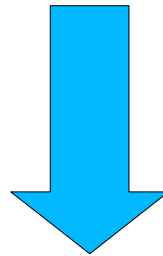
$$p(\underline{r}; t) = \left(\begin{array}{c} \text{particle moves} \\ \text{first time} \end{array} \right) + \left(\begin{array}{c} \text{particle already} \\ \text{moved} \end{array} \right)$$



Diffusion as continuous time random walk

probability of
jump with
length \underline{r} and
time step t

$$p(\underline{r}; t) = \int_{R_D} d\underline{r}' \int_t^\infty dt' \delta(\underline{r} - \underline{r}') \psi(\underline{r}', t')$$
$$+ \int_{R_D} d\underline{r}' \int_0^t dt' p(t - t'; \underline{r} - \underline{r}') \psi(\underline{r}', t')$$

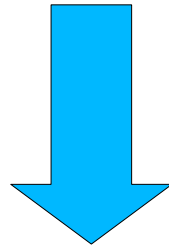


$$p(\underline{k}; u) = \frac{1 - \psi(\underline{k}, u)}{u} + p(\underline{k}; u) \psi(\underline{k}, u)$$

Diffusion as continuous time random walk

$$u p(\underline{k}; u) - 1 = K(\underline{k}; u) p(\underline{k}; u)$$

$$p(\underline{k}; u) = \frac{1 - \psi(u)}{u} + p(\underline{k}; u) \psi(\underline{k}, u)$$

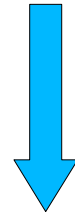


$$K(\underline{k}; u) = u \frac{\psi(\underline{k}; u) - \psi(u)}{1 - \psi(u)}$$

Diffusion as continuous time random walk

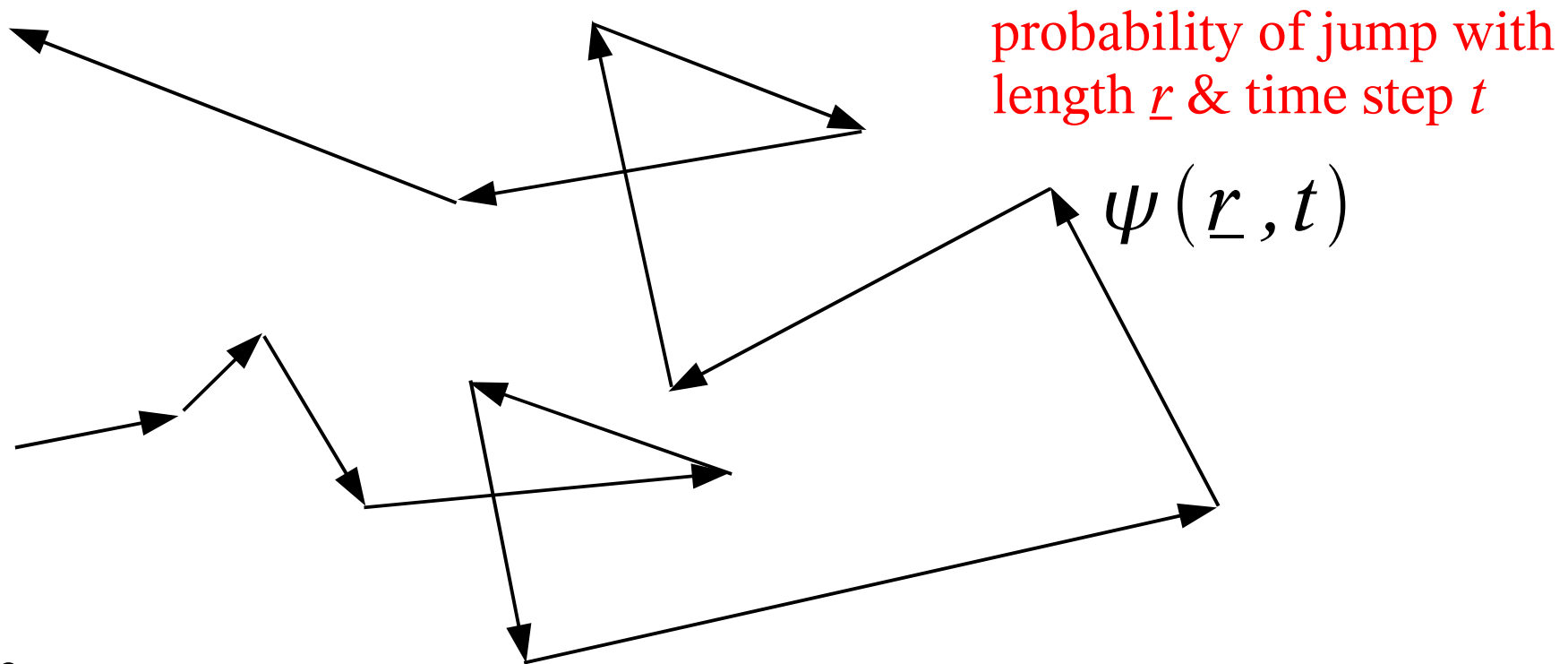
$$K(\underline{k}; u) = u \frac{\psi(\underline{k}; u) - \psi(u)}{1 - \psi(u)}$$

assume decoupled
time and length



$$K(\underline{k}; u) = (p(\underline{k}) - 1) \frac{u \psi(u)}{1 - \psi(u)}$$

Diffusion as continuous time random walk



$$\tau = \int_0^{\infty} dt t \int_{R_D} d\underline{r} \psi(\underline{r}, t) \quad \text{waiting time}$$

$$\sigma^2 = \int_0^{\infty} dt \int_{R_D} d\underline{r} r^2 \psi(\underline{r}, t) \quad \text{space step}$$

Diffusion as continuous time random walk

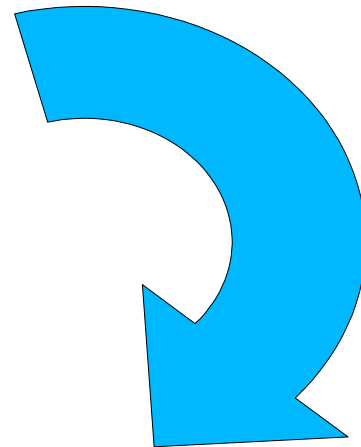
$$\tau = \int_0^{\infty} dt t \int_{R_D} d\underline{r} \psi(\underline{r}, t)$$

$$t \rightarrow \infty,$$

$$k \rightarrow 0,$$

$$u \rightarrow 0$$

$$\sigma^2 = \int_0^{\infty} dt \int_{R_D} d\underline{r} r^2 \psi(\underline{r}, t)$$



$$\tau < \infty, \sigma^2 < \infty \quad \phi(u) = 1 - \tau u, \quad p(k) = 1 - \frac{\sigma^2}{2} k^2$$

$$\tau = \infty, \sigma^2 < \infty \quad \phi(u) = 1 - c_1 u^\alpha \quad (0 < \alpha < 1),$$

$$p(k) = 1 - \frac{\sigma^2}{2} k^2$$

Diffusion as continuous time random walk

$$\tau < \infty, \sigma^2 < \infty \quad \phi(u) = 1 - \tau u, \quad p(k) = 1 - \frac{\sigma^2}{2} k^2$$

$$up(\underline{k}; u) - 1 = \frac{\sigma^2}{2\tau} k^2 p(\underline{k}; u) \quad \text{normal diffusion}$$

$$\frac{\partial}{\partial t} p(\underline{r}; t) = \frac{\sigma^2}{2\tau} \frac{\partial^2}{\partial r^2} p(\underline{r}, t')$$

$$\boxed{\tau = \infty}, \sigma^2 < \infty \quad \phi(u) = 1 - c_1 u^\alpha \quad (0 < \alpha < 1),$$

$$\text{anomalous diffusion} \quad p(k) = 1 - \frac{\sigma^2}{2} k^2$$

$$\frac{\partial}{\partial t} p(\underline{r}; t) = \frac{\sigma^2}{2\tau} \boxed{\frac{\partial^\alpha}{\partial t^\alpha}} \frac{\partial^2}{\partial r^2} p(\underline{r}, t')$$

Types of diffusion

$$\tau < \infty, \sigma^2 < \infty$$

normal diffusion

$$\langle r^2(t) \rangle = \int_{R_D} d\underline{r} r^2 p(\underline{r}, t) = 6Dt$$

$$\tau = \infty, \sigma^2 < \infty$$

sub-diffusion

$$\langle r^2(t) \rangle = \int_{R_D} d\underline{r} r^2 p(\underline{r}, t) \sim t^\alpha, \quad 0 < \alpha < 1$$

time and length have
to be coupled

super-diffusion

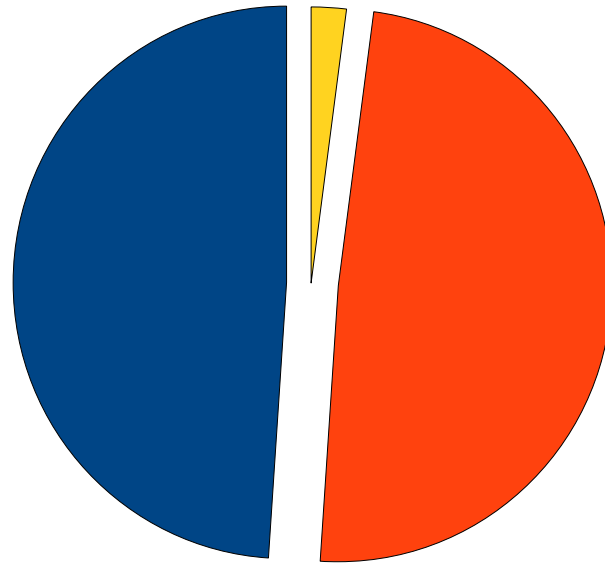
$$\langle r^2(t) \rangle = \int_{R_D} d\underline{r} r^2 p(\underline{r}, t) \sim t^\alpha, \quad \alpha > 1$$

How frequent are different types of the diffusion?

**sub-diffusion
(limiting case)**

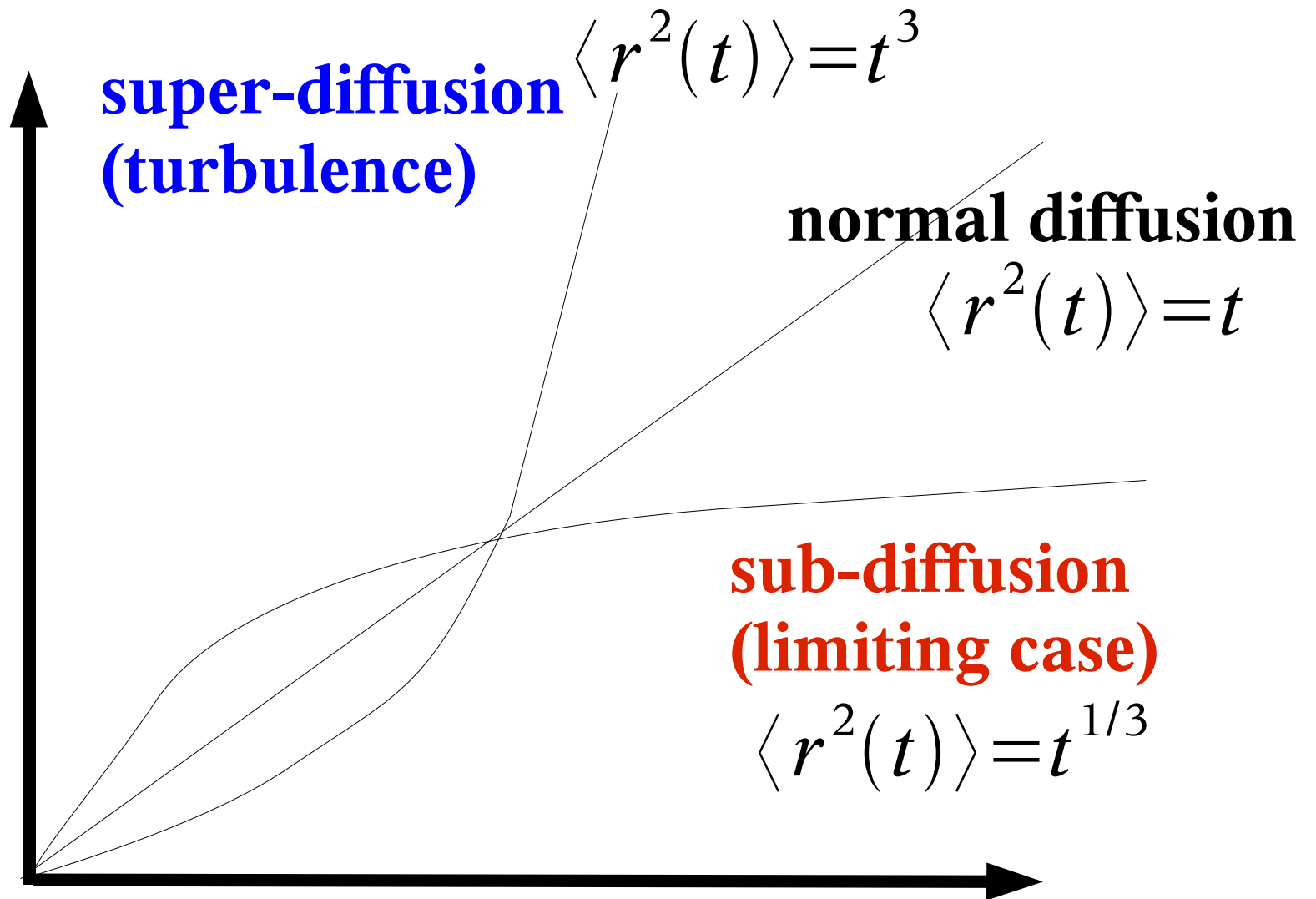
$$\tau = \infty, \sigma^2 < \infty$$

normal diffusion
 $\tau < \infty, \sigma^2 < \infty$

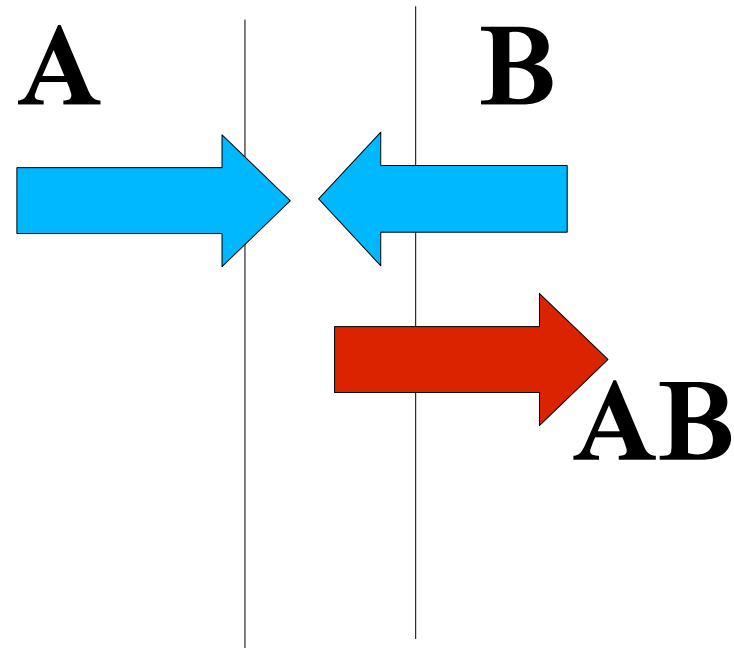


**super-diffusion
(turbulence)**

$$\langle r^2(t) \rangle = t^3$$

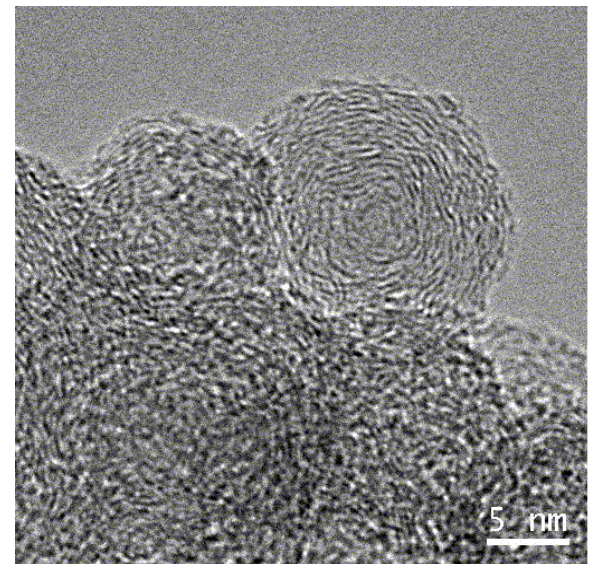
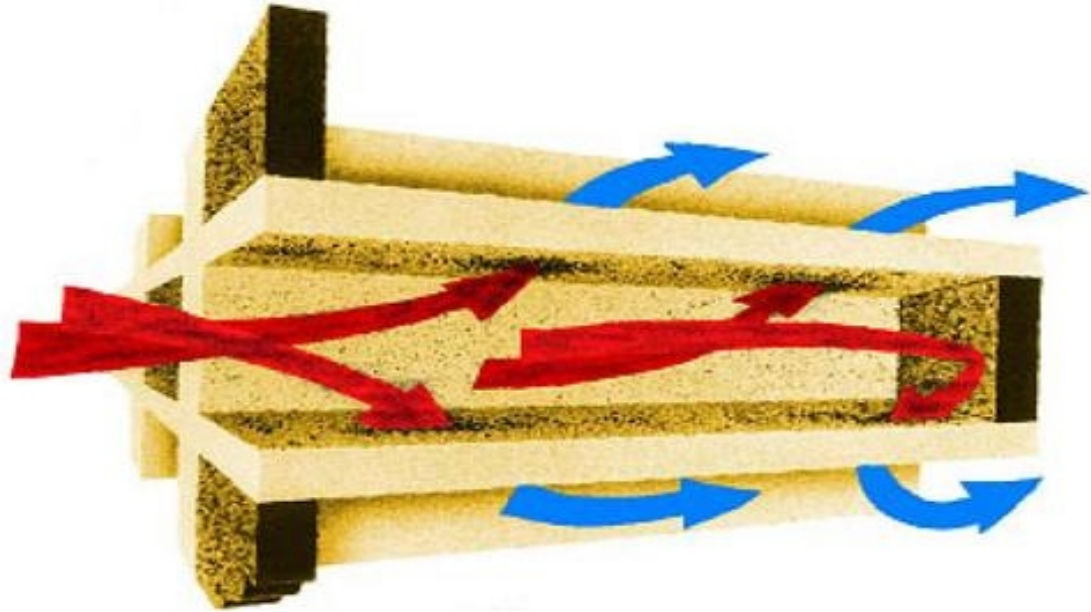
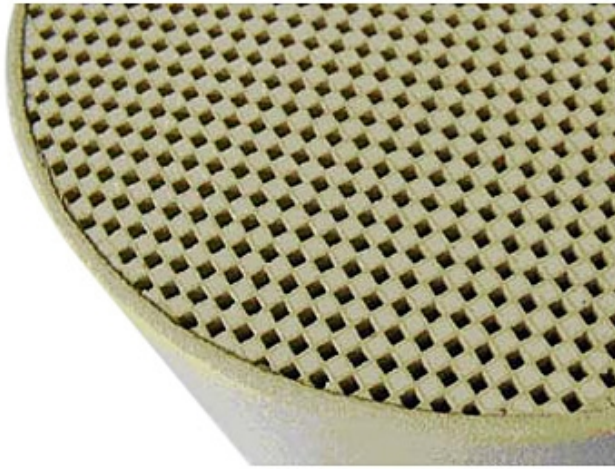


Some applications:



selective membranes

Some applications:



Creation of the model porous structures



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Total energy in the system is

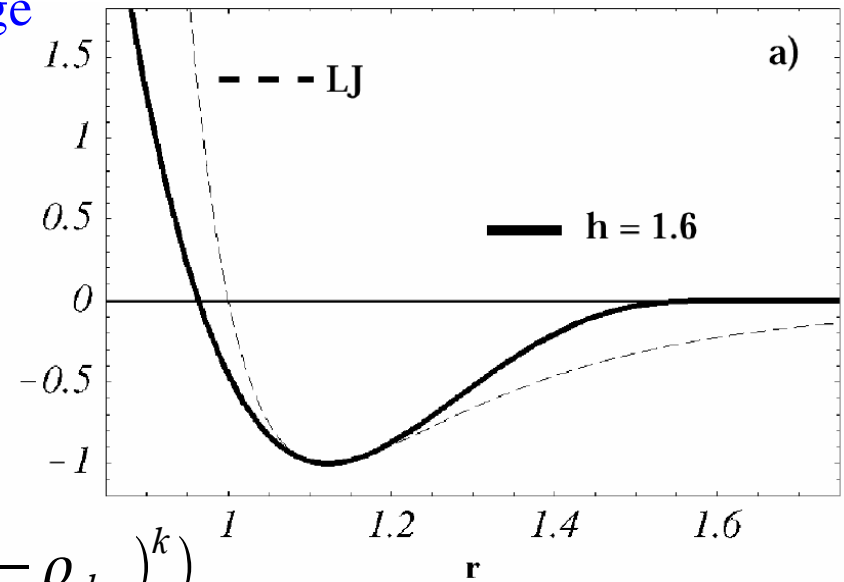
expressed as sum of two contributions:

$$E_{tot} = \sum_i \left(\sum_j \Phi(r_{ij}) + F(\rho_i) \right)$$

$$\Phi(r) = \phi_0 r_0^{-6} \left(3(r - r_{cut})^4 - 4(r_{cut} - r_{min})(r_{cut} - r)^3 \right), r \leq r_{cut}$$

$\Phi(r) = 0, r > r_{cut}$ binary contribution - short range attractive potential (SHRAT)

$$r_{min} = 2^{1/6}, r_{cut} = 1.6, w_0 = \frac{105}{16 \pi r_{cut}^3}, \rho_{des} = 1$$



$$F(\rho_i) = \phi_0 \sum_{k=2,4,\dots} r_0^{3k} F_k \left((\rho_i - \rho_{des})^k - (w(0) - \rho_{des})^k \right)$$



many-body contribution embedding functional

$$\rho_i = \sum_j w(r_{ij})$$

$$w(r) = w_0 \left(1 + 3 \frac{r}{r_{cut}} \right) \left(1 - \frac{r}{r_{cut}} \right)^3$$

Lucy's weight function



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<http://www.fisica.uniud.it/~ercolessi/md/>

Molecular dynamics

method: Basics

$$\dot{\mathbf{r}}_i = \mathbf{v}_i$$

$$\dot{\mathbf{v}}_i = -\nabla_i V(\mathbf{r}_1, \dots, \mathbf{r}_N)$$

molecules, atoms

metals, semiconductors

$$V(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N) = \sum_i \sum_{j>i} \phi(|\mathbf{r}_i - \mathbf{r}_j|)$$

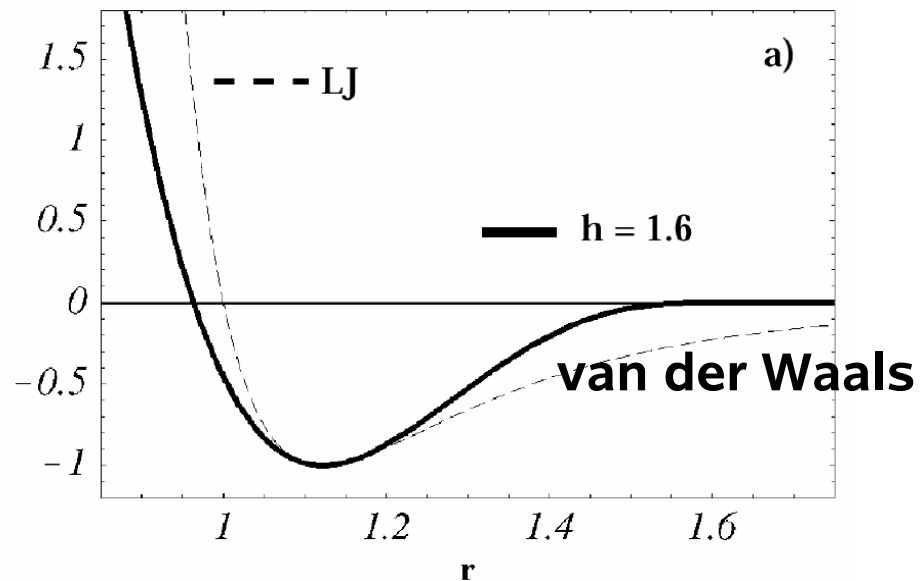
$$V(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N) = \sum_i \left(\sum_{j>i} \phi(|\mathbf{r}_i - \mathbf{r}_j|) + \Theta(\rho_i) \right)$$

$$\Phi(r) = \phi_0 r_0^{-6} (3(r - r_{cut})^4 - 4(r_{cut} - r_{min})(r_{cut} - r)^3), r \leq r_{cut}$$

$$\Phi(r) = 0, r > r_{cut}$$

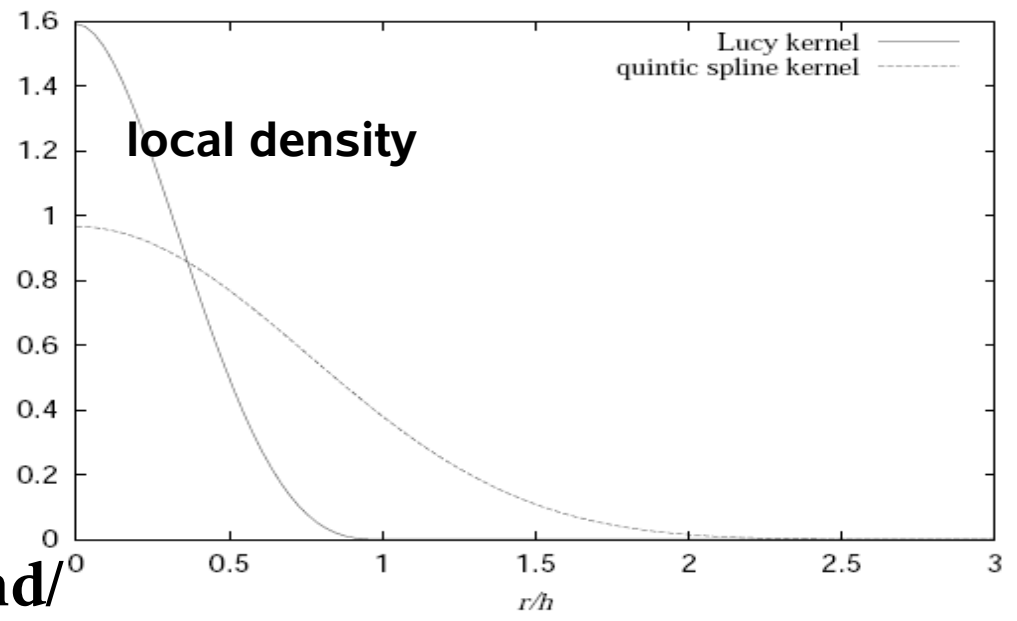
$$\Theta(\rho_i) = \phi_0 \sum_{k=2,4,\dots} r_0^3 k F_k ((\rho_i - \rho_{des})^k - (w(0) - \rho_{des})^k)$$

Paulie principle



$$\rho_i = \sum_j w(r_{ij})$$

$$w(r) = w_0 \left(1 + 3 \frac{r}{r_{cut}}\right) \left(1 - \frac{r}{r_{cut}}\right)^3$$



Molecular dynamics method: What can we get out of it?

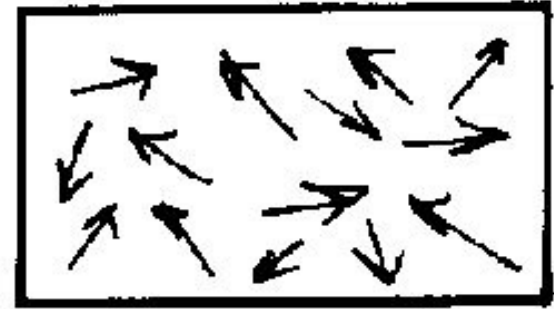
$$V(t) = \sum_i \sum_{j>i} \phi(|\mathbf{r}_i(t) - \mathbf{r}_j(t)|) \quad \text{potential energy}$$

$$K(t) = \frac{1}{2} \sum_i m_i [v_i(t)]^2 \quad \text{kinetic energy}$$

$$E(t) = K(t) + V(t) \quad \text{total energy}$$

$$K(t) = \frac{3}{2} N k_B T(t) \quad \text{temperature from kinetic energy}$$

$$PV = N k_B T - \frac{1}{3} \left\langle \sum_i \mathbf{r}_i \cdot \mathbf{F}_i \right\rangle \quad \text{pressure}$$

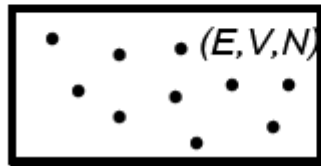


$\{q_i, p_i\}$

Molecular dynamics method: How to keep temperature constant?

- Microcanonical ensemble

isolated system



*energy transfer not allowed
E constant*

$$f(\Gamma) = \delta(H(\Gamma) - E)$$

- Canonical ensemble

heat reservoir T



*energy transfer allowed
T constant*

$$f(\Gamma) = \exp(-H(\Gamma)/k_B T)$$

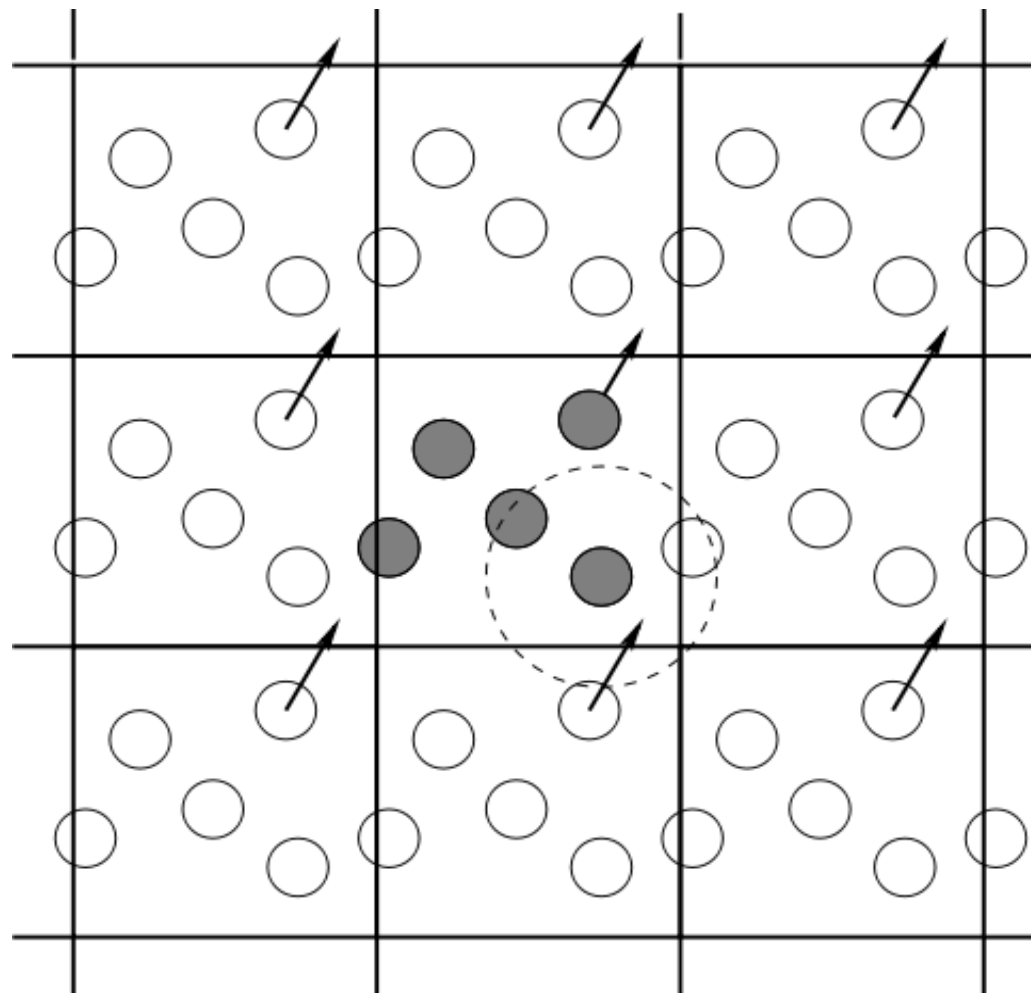


$$\hat{\mathbf{v}}_i \left(t + \frac{\Delta t}{2} \right) = S \mathbf{v}_i \left(t + \frac{\Delta t}{2} \right)$$



$$S = \left[\frac{\frac{3}{2} N k_B T}{\sum_i \frac{m_i}{2} \left[\mathbf{v}_i \left(t + \frac{\Delta t}{2} \right) \right]^2} \right]^{1/2}$$

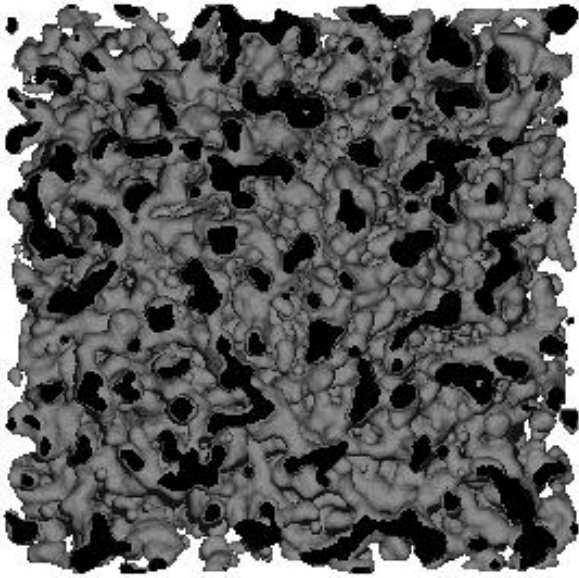
Molecular dynamics method: How to simulate bulk material?



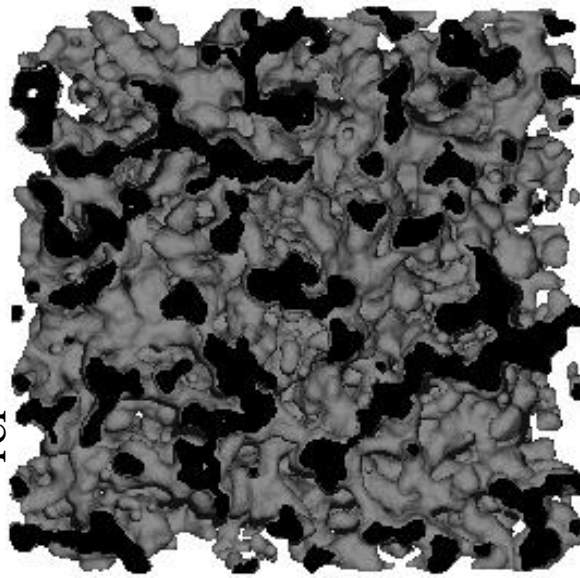
periodic boundary conditions

<http://www.fisica.uniud.it/~ercolessi/md/>

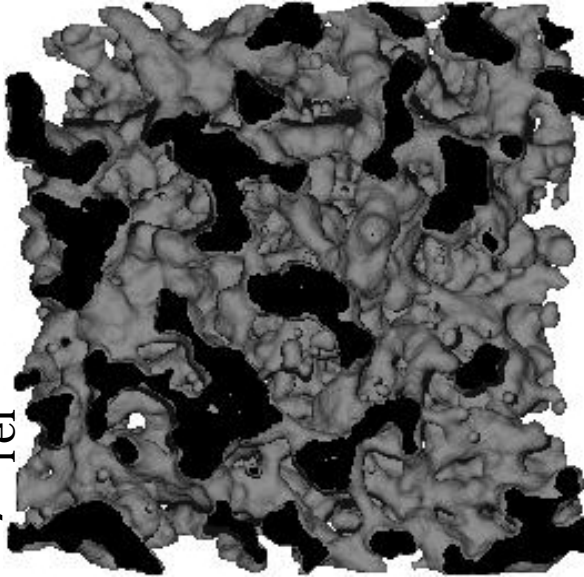
$t/t_{\text{ref}}=100$



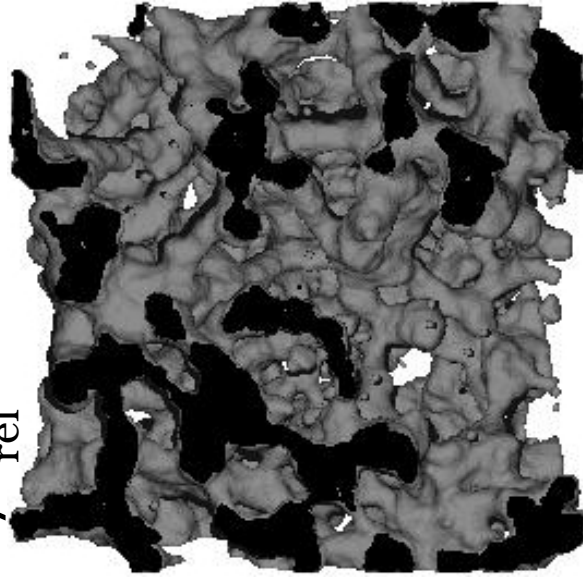
$t/t_{\text{ref}}=200$



$t/t_{\text{ref}}=500$

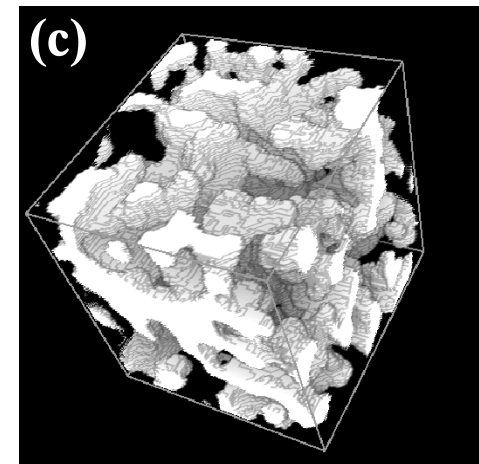
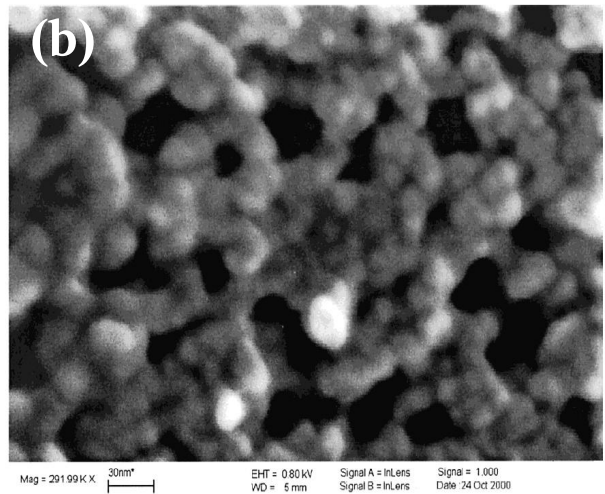
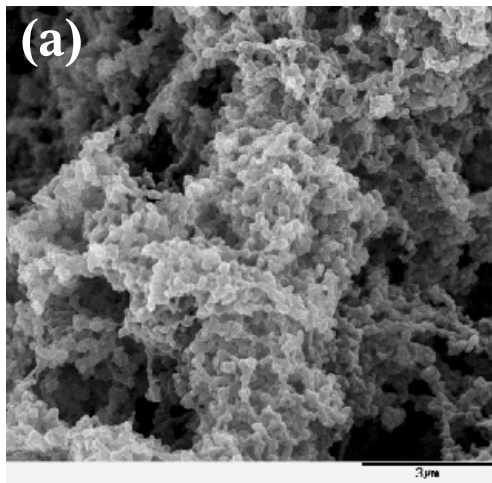


$t/t_{\text{ref}}=2000$



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- (a) U.Mähr, H. Purnama, E. Kempin, R. Schomäcker, and K.-H. Reichert, *J. Membr. Sci.*, **171** (2000), 285.
(b) M. Meyer, A. Fischer, and H. Hoffman, *J. Phys. Chem. B*, **106** (2002) 1528.
(c) <http://ciks.cbt.nist.gov/~garbocz/paper75/paper75.html>

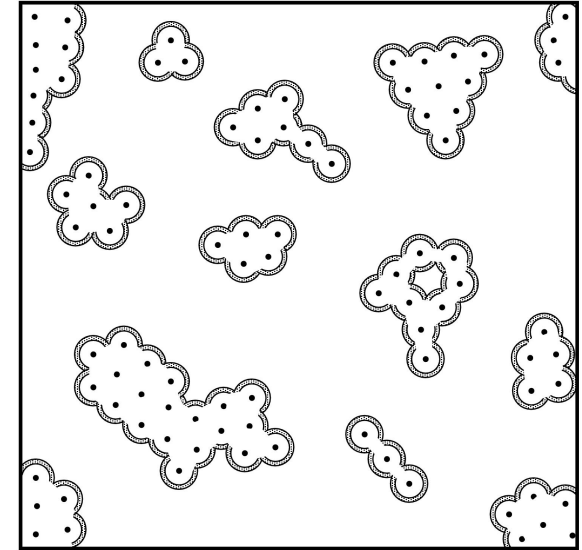


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Monte Carlo integration:

volume - points from a random sample are counted, if they are within certain radius around the particles (r): $V/V_0 = N_{\text{within}}/N_{\text{tot}}$.
surface is calculated as numerical differential of volume over radius: $\partial V(r)/\partial r$.



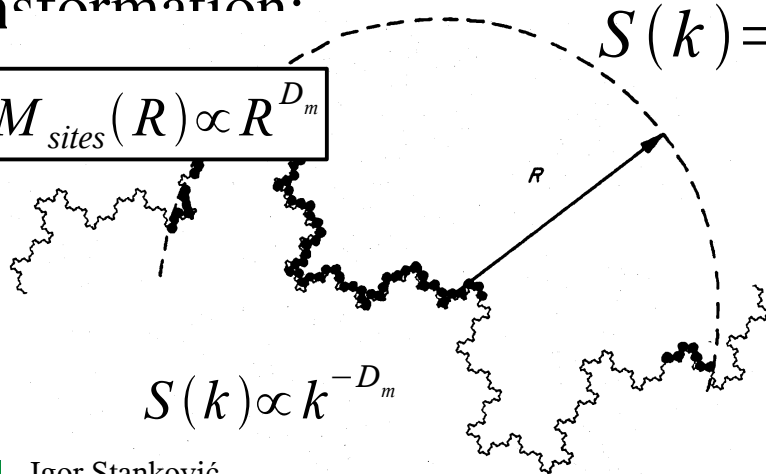
Structure factor and fractal dimension:

structure factor $S(k)$ is related to the *pair distribution function*, $g(r)$, by the Fourier transformation:

$$S(k) = 1 + 4\pi n \int_0^{\infty} (g(r) - 1) r^2 \frac{\sin(kr)}{kr} dr$$

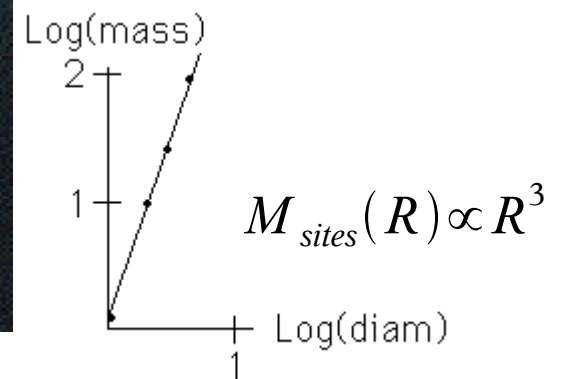
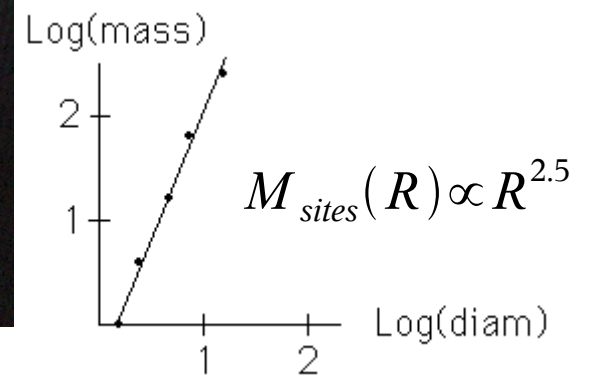
$$M_{\text{sites}}(R) \propto R^{D_m}$$

$$S(k) \propto k^{-D_m}$$



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<http://classes.yale.edu/fractals/FracAndDim/BoxDim/PowerLaw/PowerLaw.html>

Mass and surface fractals:

$$M_{\text{sites}}(R) \propto R^{D_m}$$

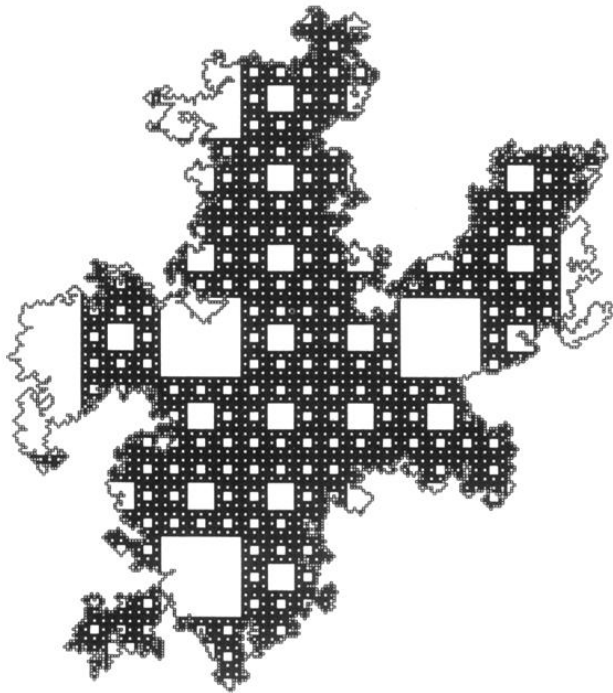
$$M_{\text{sites on the surface}}(R) \propto R^{D_s} \quad \text{for small } r!$$

$$g(r) \propto r^{D_m - 3}$$

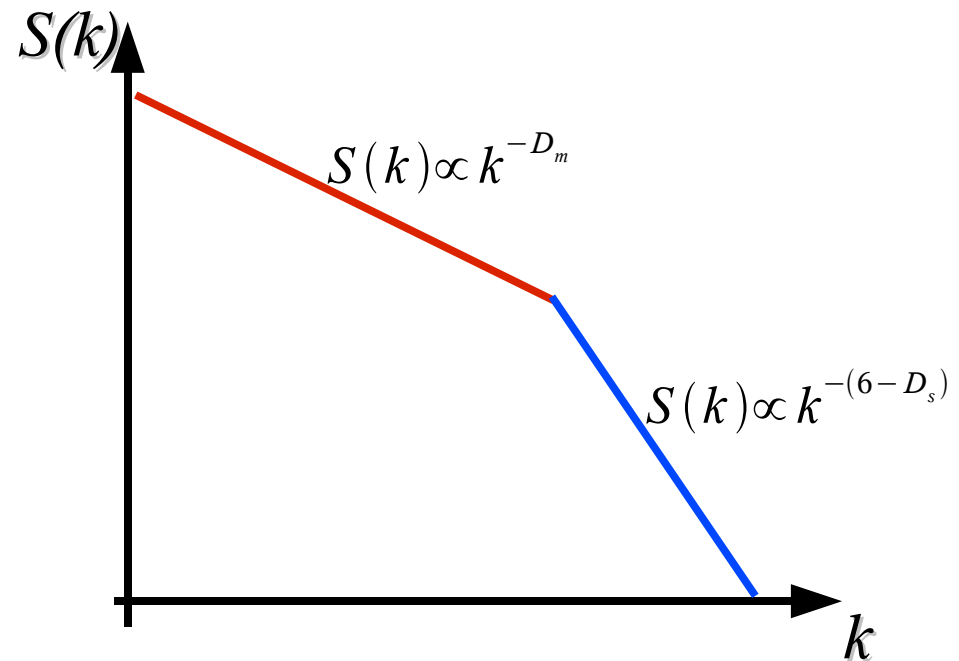
$$g(r) \propto 1 - r^{3 - D_s}$$

$$S(k) \propto k^{-D_m}$$

$$S(k) \propto k^{-(6 - D_s)}$$



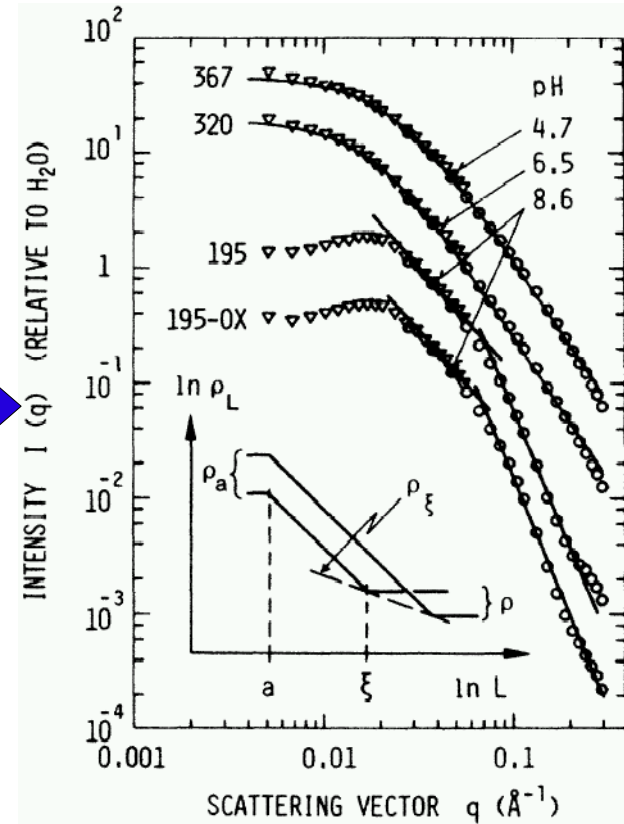
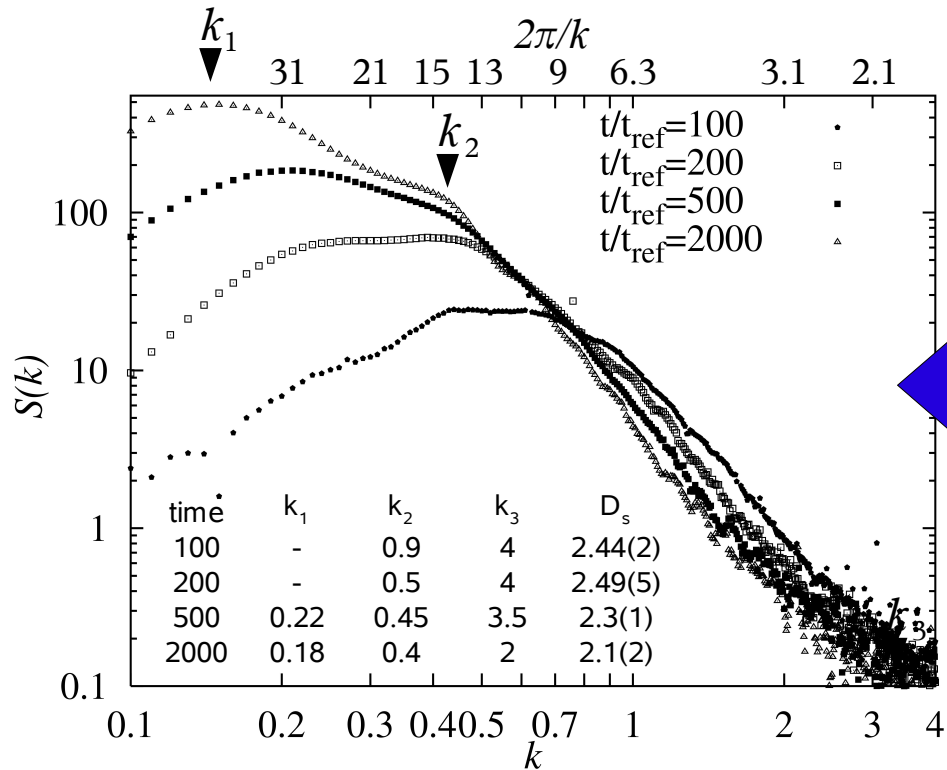
in reality structures are fractal only on a certain range of scales:



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GEAM ($F_2=1, F_k=0, k>2$),
 $n=0.25, T=0.01$



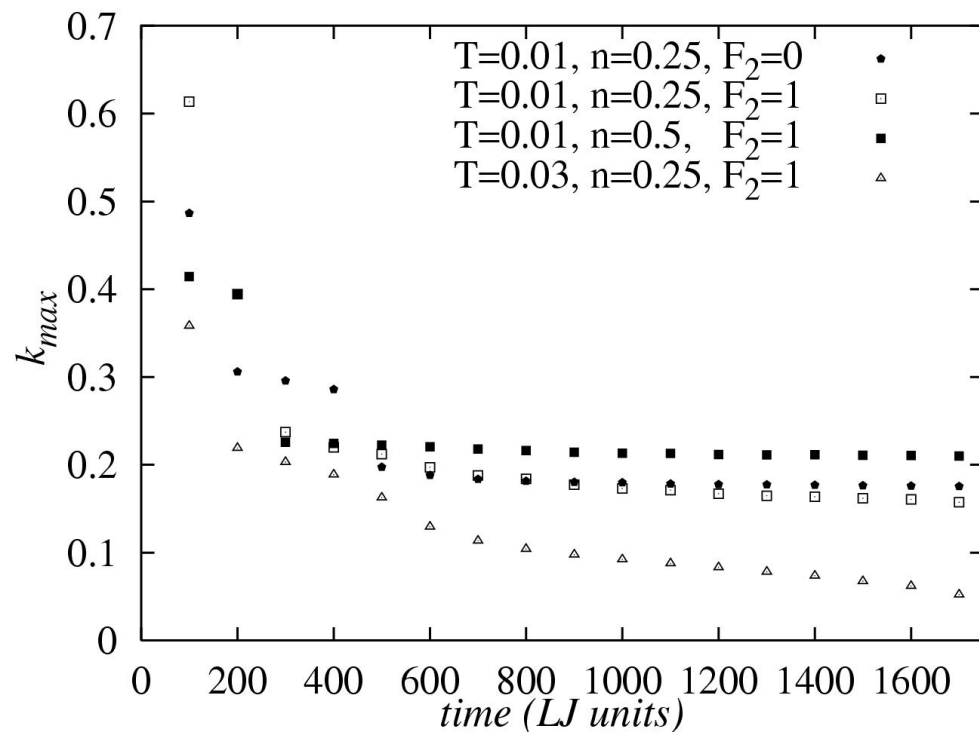
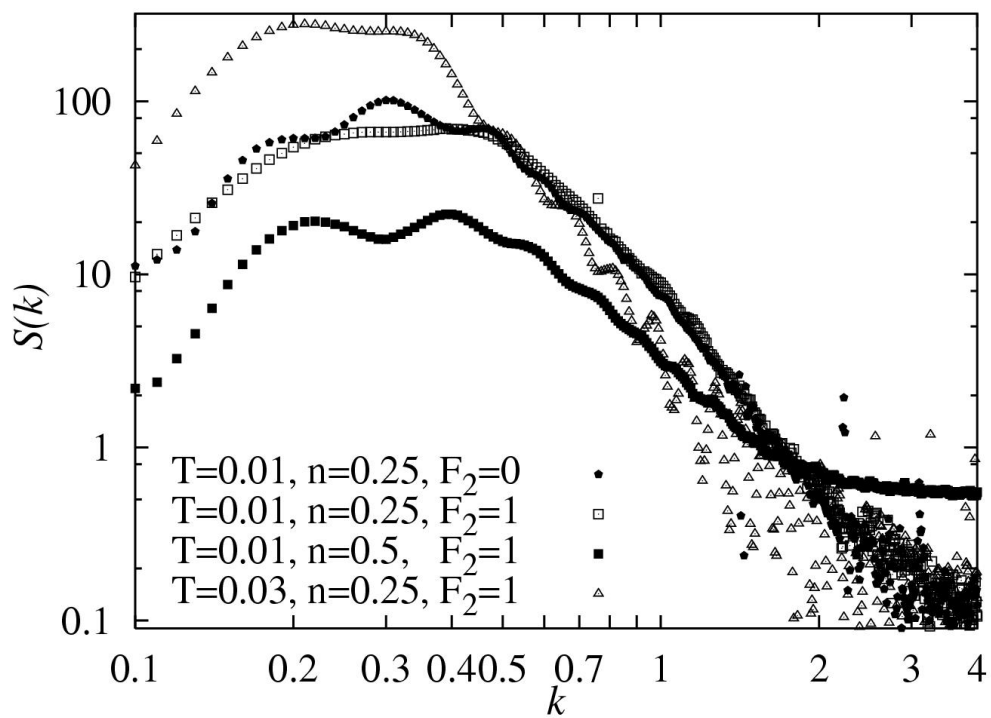
$D_s=2.4$

R. Vacher et al., Phys. Rev. B, **37** (1988), 6500.



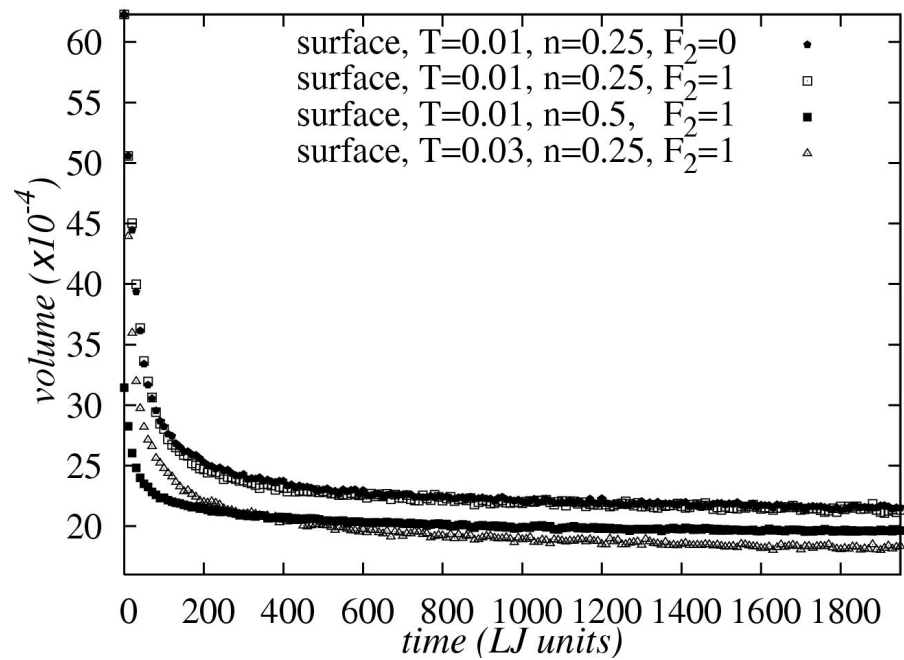
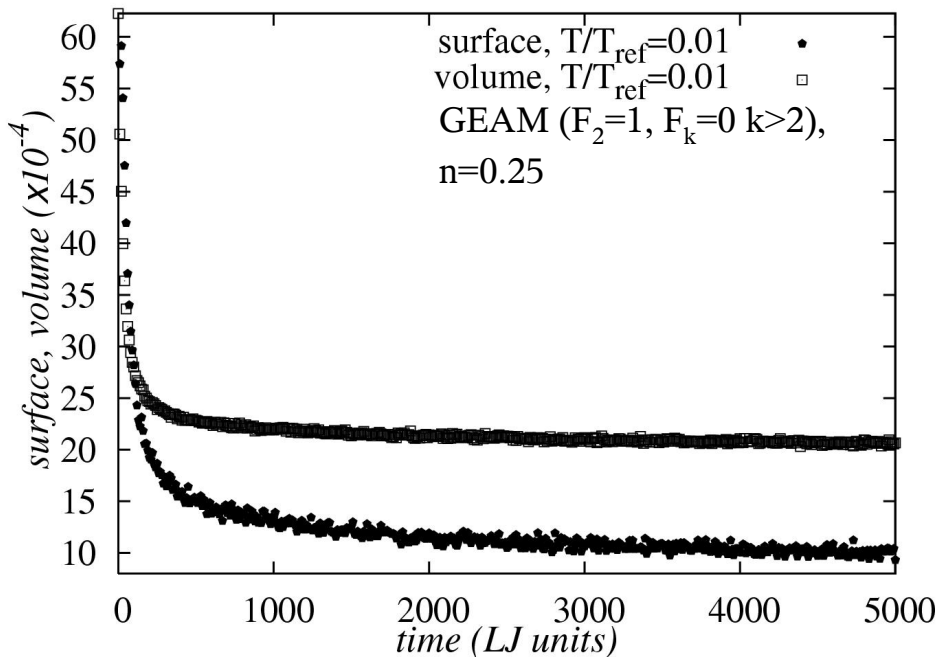
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T	n	F_2	D
0.01	0.25	0	$D_m=2.5(2)$
0.01	0.25	1	$D_s=2.7(4)$
0.01	0.5	1	$D_s=2.5(2)$
0.03	0.25	1	$D_s=2.2(2)$

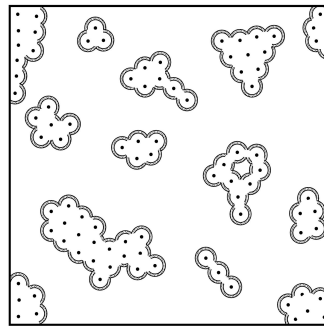
Volume and surface:



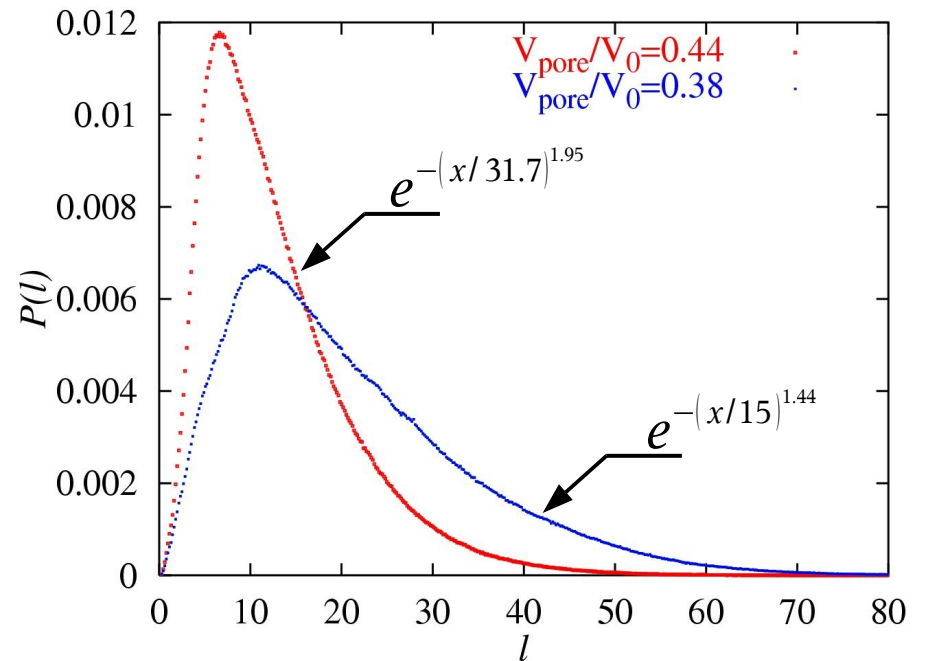
Monte Carlo integration:

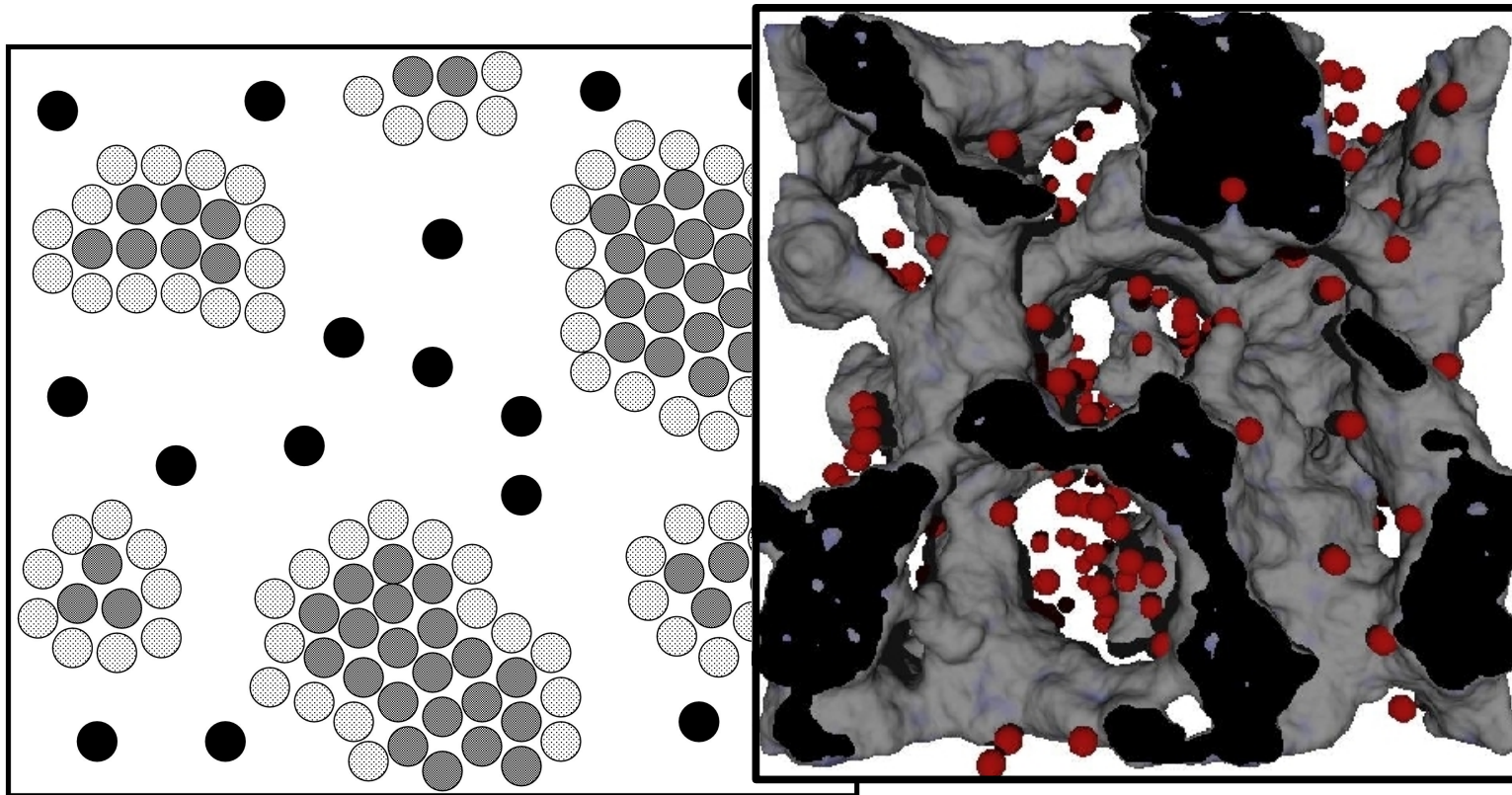
volume - $V/V_0 = N/N_{tot}$,

surface - $\partial V(r)/\partial r$.



distribution of free flight paths





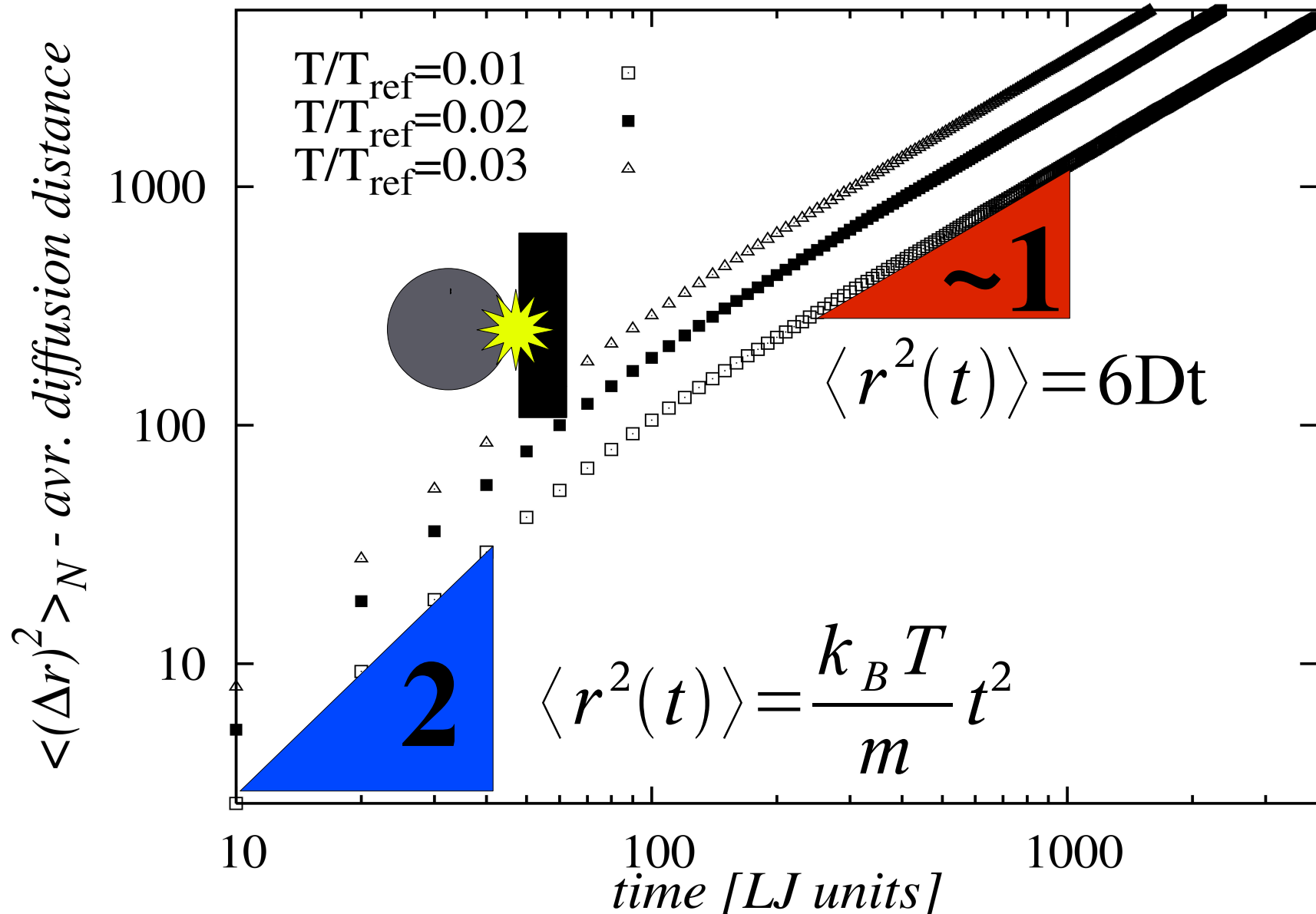
- gas particles
- particles with fixed positions inside of the sponge
- sponge wall particles

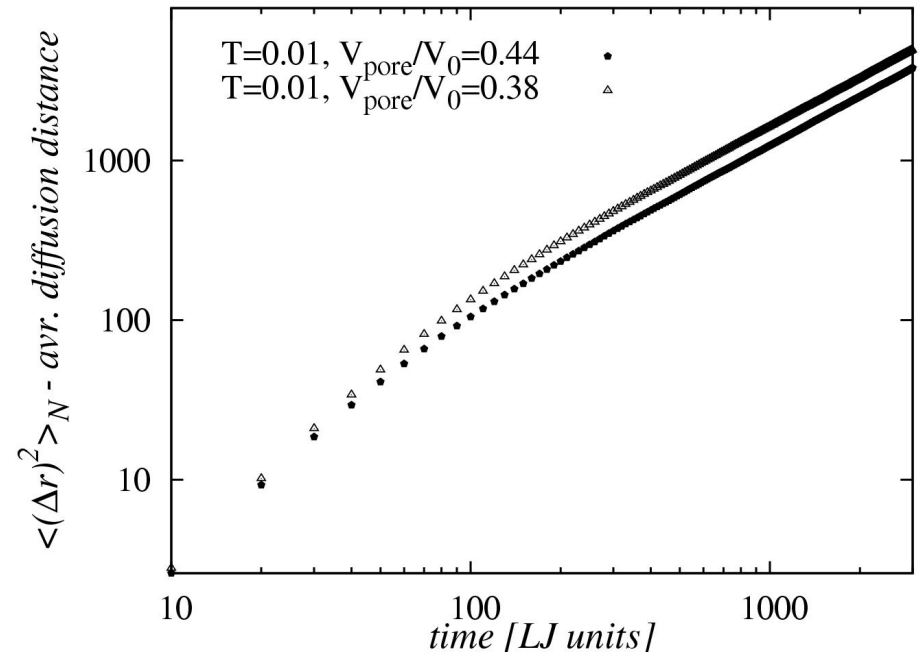
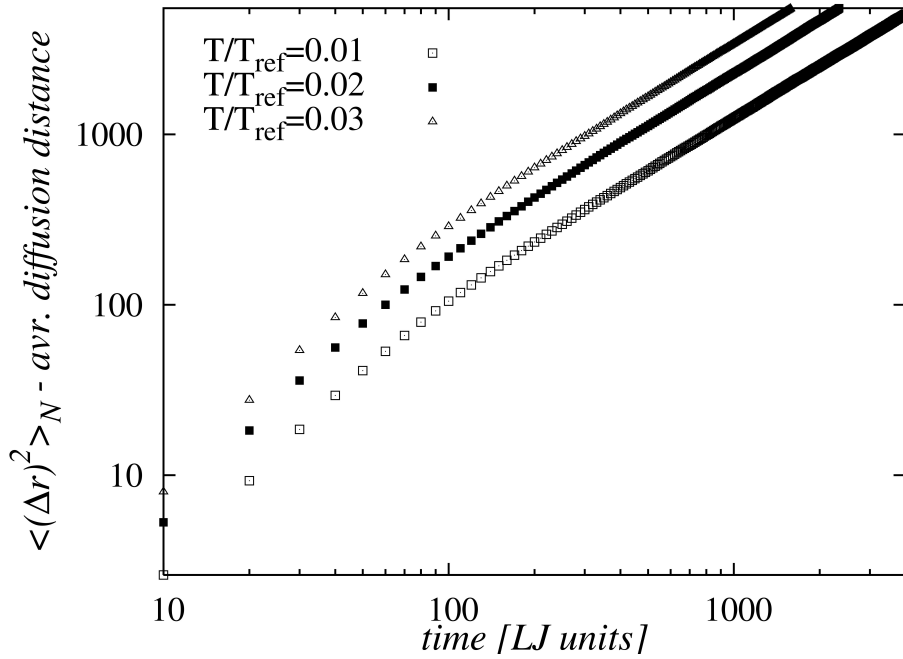
- gas particles do not interact
- gas-wall interaction SHREP
- thermostat acts only on wall particles!

$$\gg \Phi(r) = \phi r_0^{-4} (r - r_{\text{cut}})^4, r \leq r_{\text{cut}}$$

$$\Phi(r) = 0, r > r_{\text{cut}} \quad (\text{repulsive})$$

Knudsen self-diffusion (small densities):



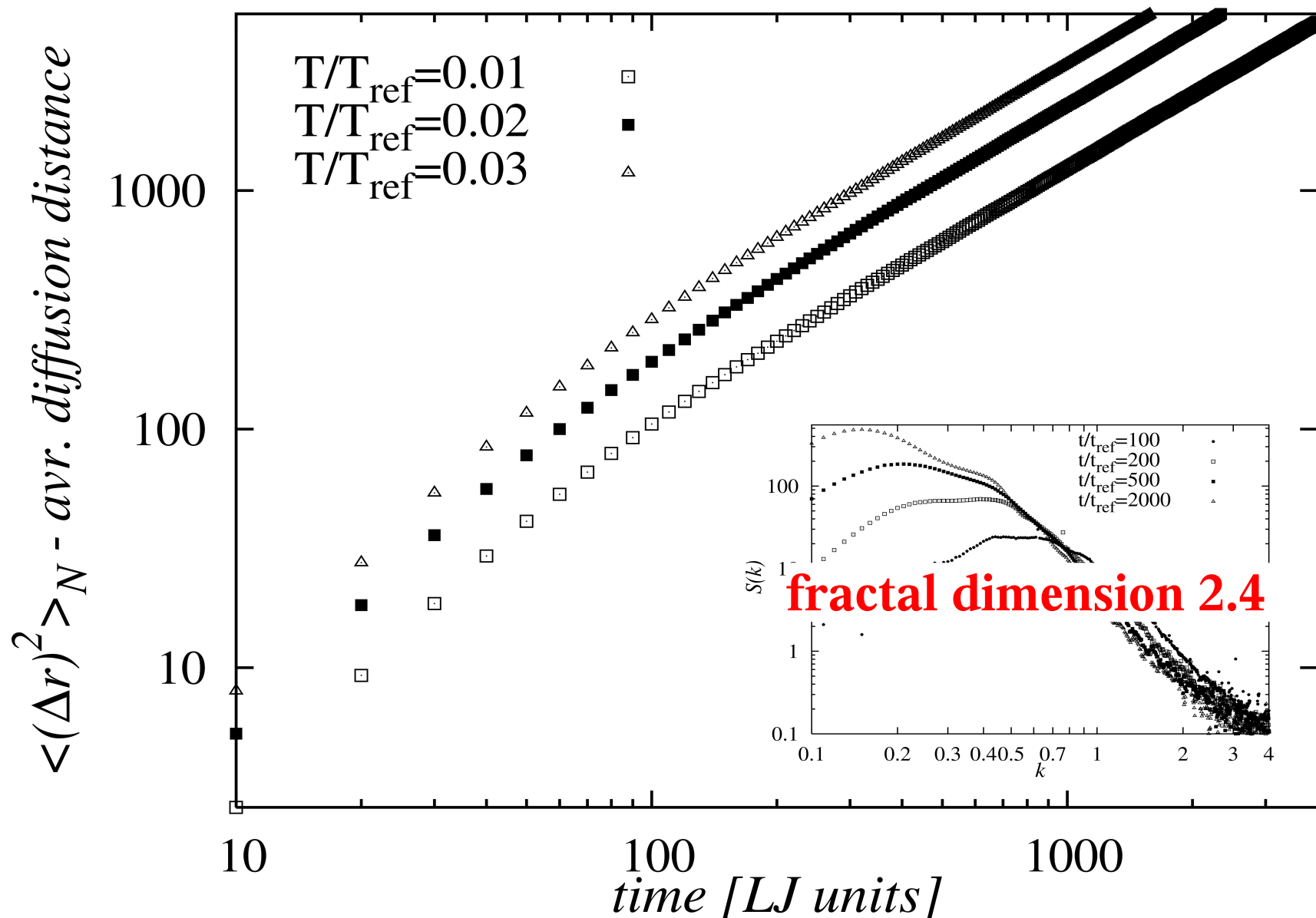


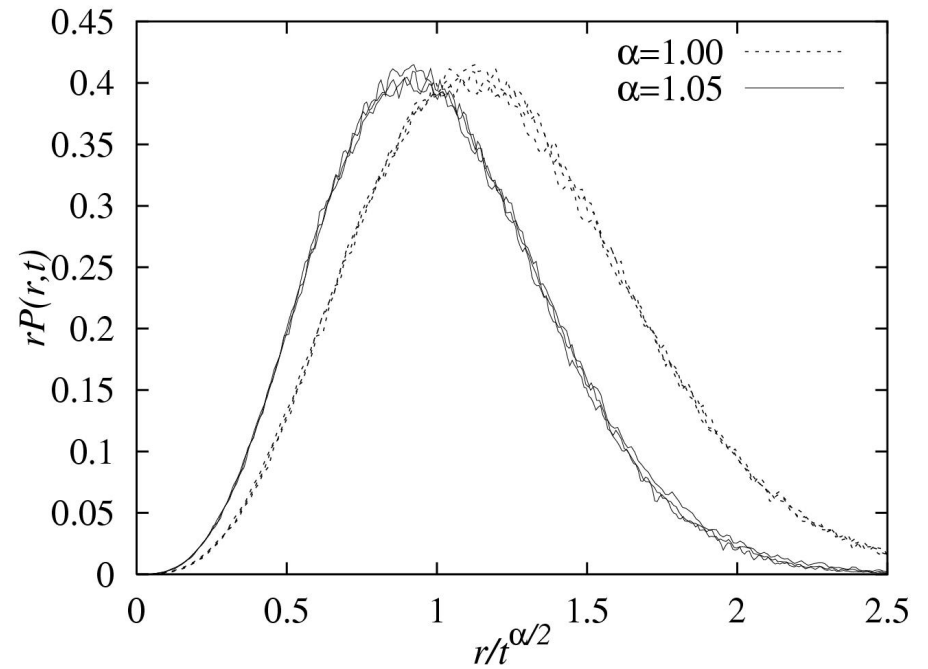
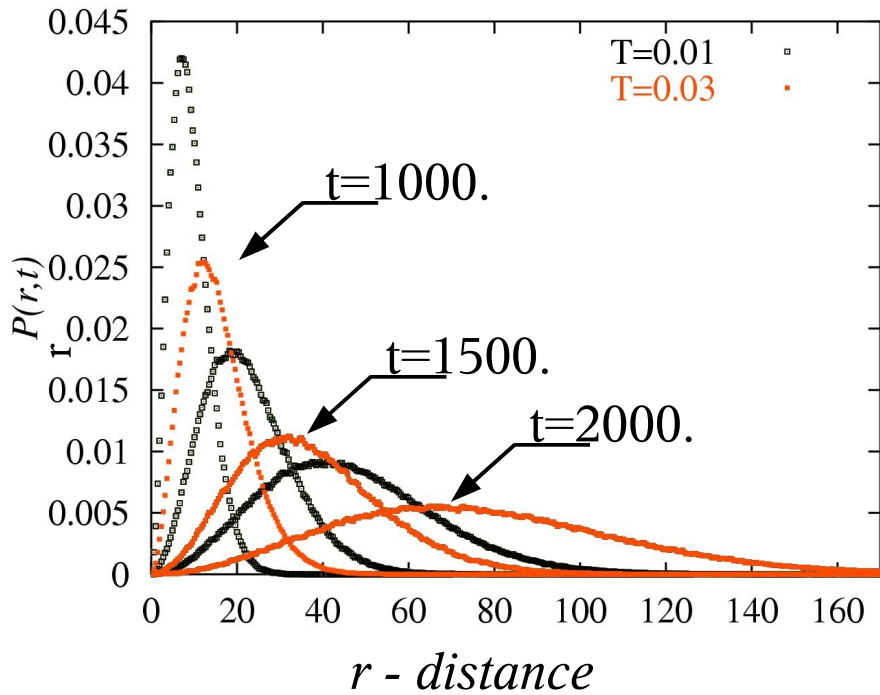
$$\langle r^2(t) \rangle \approx Dt^\alpha \quad \text{for } t/t_{\text{free}} \gg 1$$



T	V_{pore}/V_0	D	α
0.01	0.445	1.18(1)	1.008(1)
0.02	0.445	1.52(2)	1.019(2)
0.03	0.445	2.19(4)	1.055(2)
0.01	0.378	1.41(2)	0.976(2)
0.02	0.378	3.42(6)	1.058(2)
0.03	0.378	3.50(3)	1.02(1)

Knudsen self-diffusion (small densities):



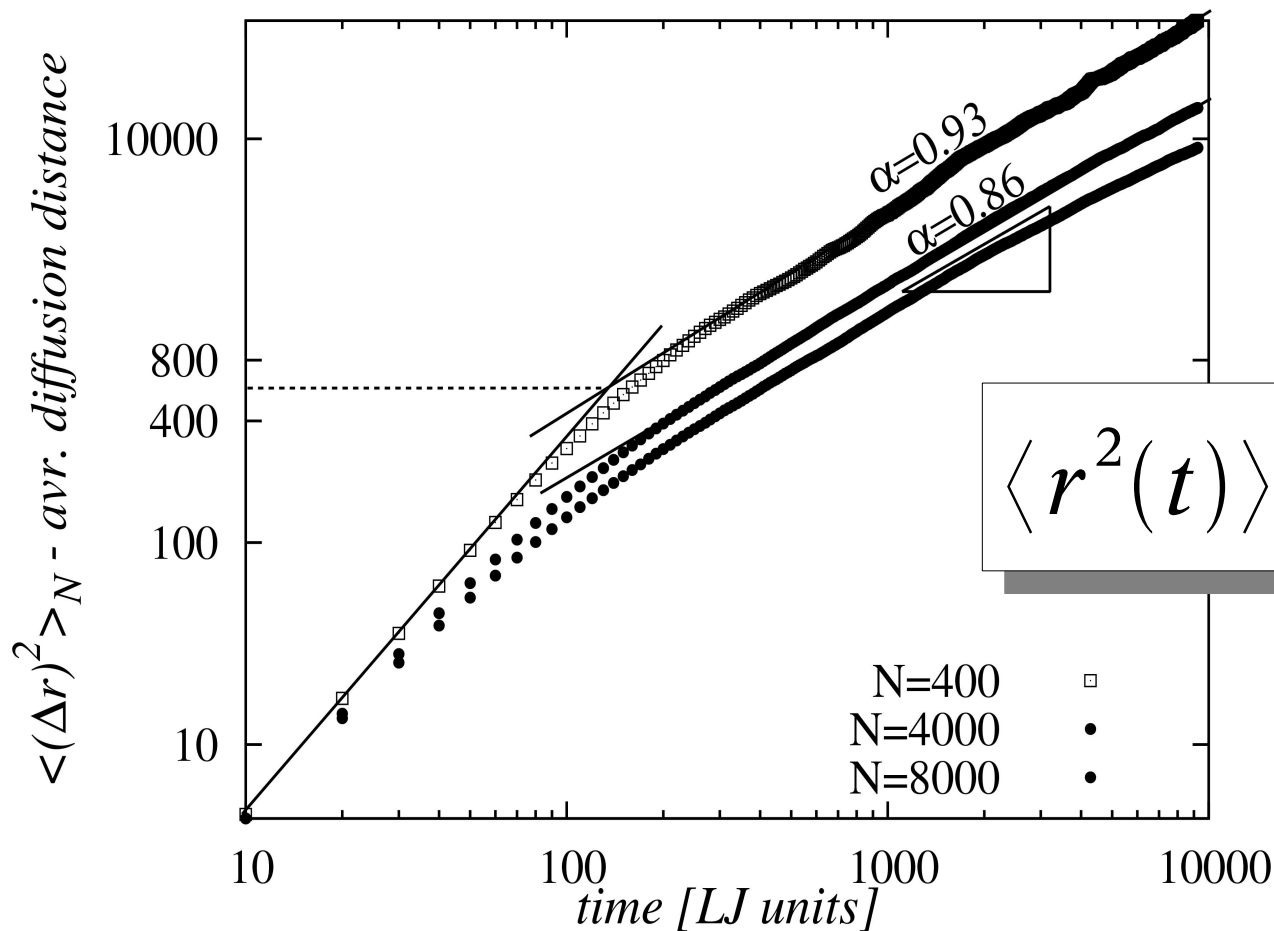
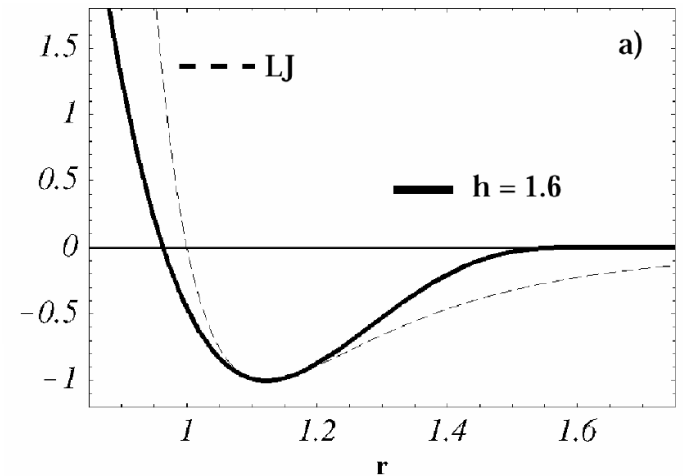


$P(r,t)$ - probability that random walker has displaced to distance r after time t

$$P(r, t) = C_{1t}^{-3\alpha/2} e^{\frac{r^2}{2Dt^\alpha}}$$

- gas-gas interaction SHRAT (attractive)

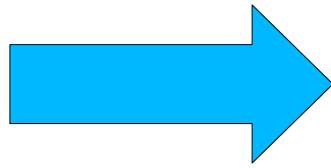
- gas-wall interaction SHREP



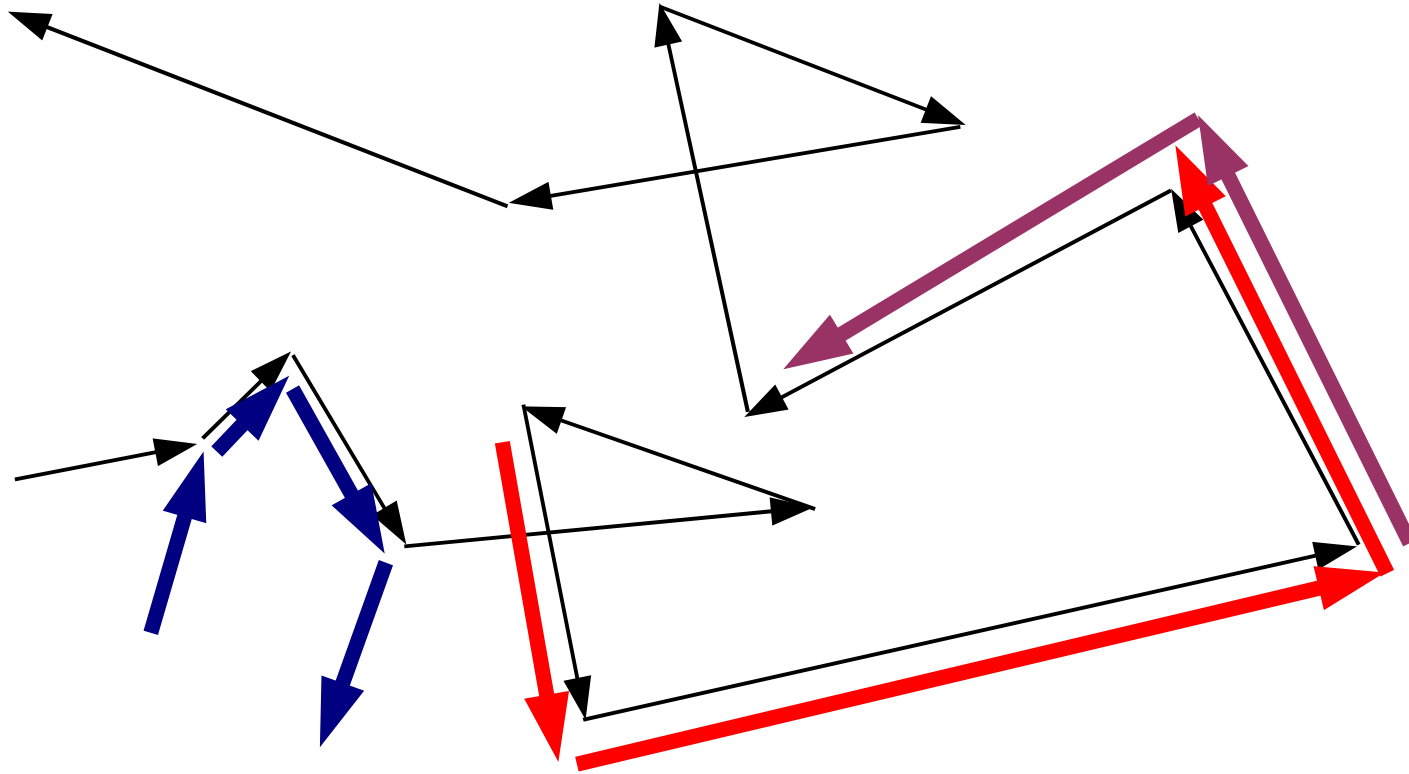
$$\langle r^2(t) \rangle \sim t^\alpha, \quad 0 < \alpha < 1$$

sub-diffusion

particles
group



mass increases/
speed decreases



$$\tau = \infty, \sigma^2 < \infty$$

as result we have
sub-diffusion